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Coupling of heat and spin currents in metallic multilayers

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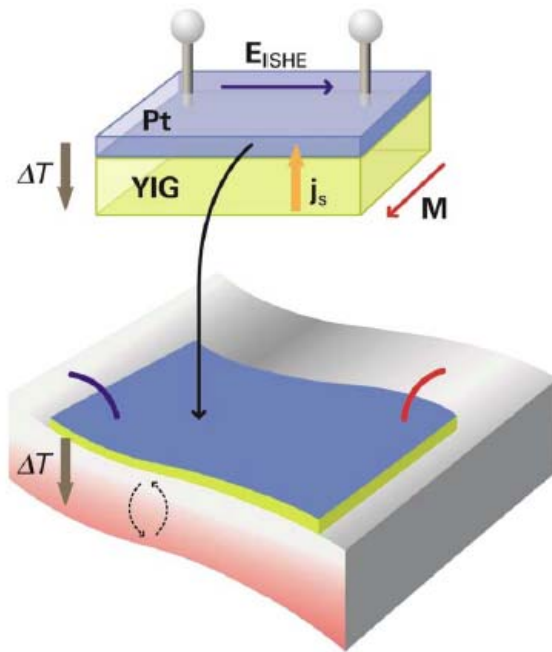
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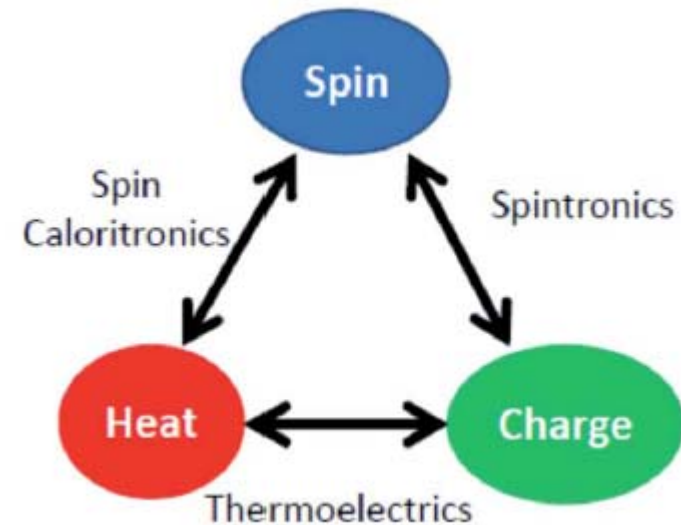
Choi *et al.*, *Nature Communications* (2014)

Motivation I: Can we make use of spin in heat engines?

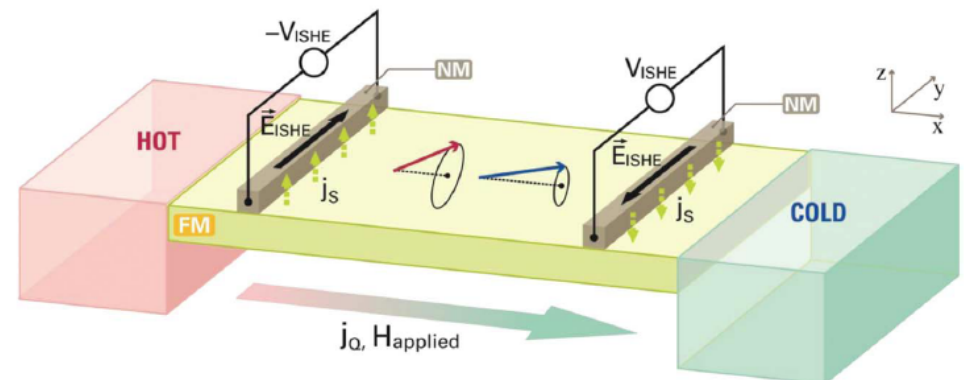
- Electronic states enumerated by energy, wave-vector, spin
- Possible advantages in geometrical scaling, $\nabla E \perp \nabla T$.



Kirihara *et al.*, *Nat. Mat.* (2013)



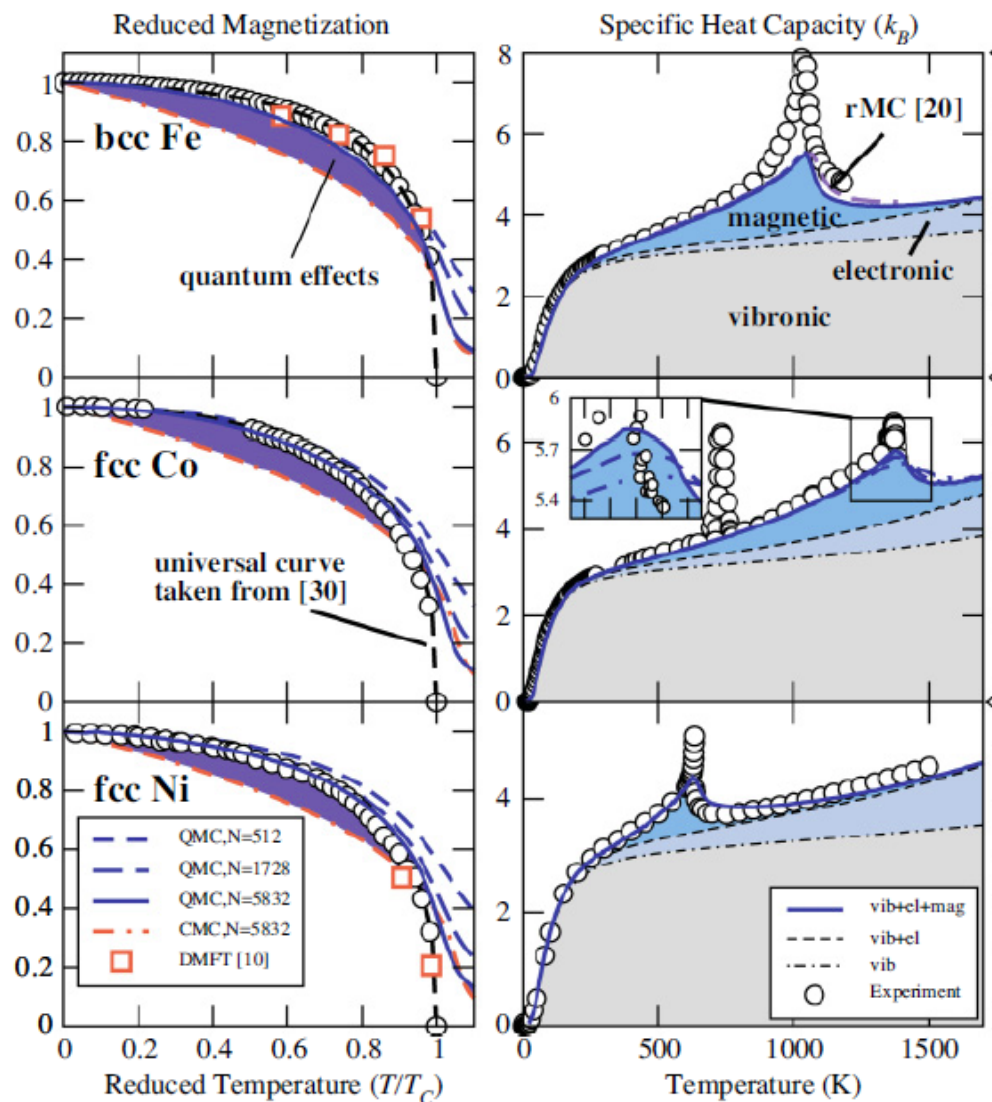
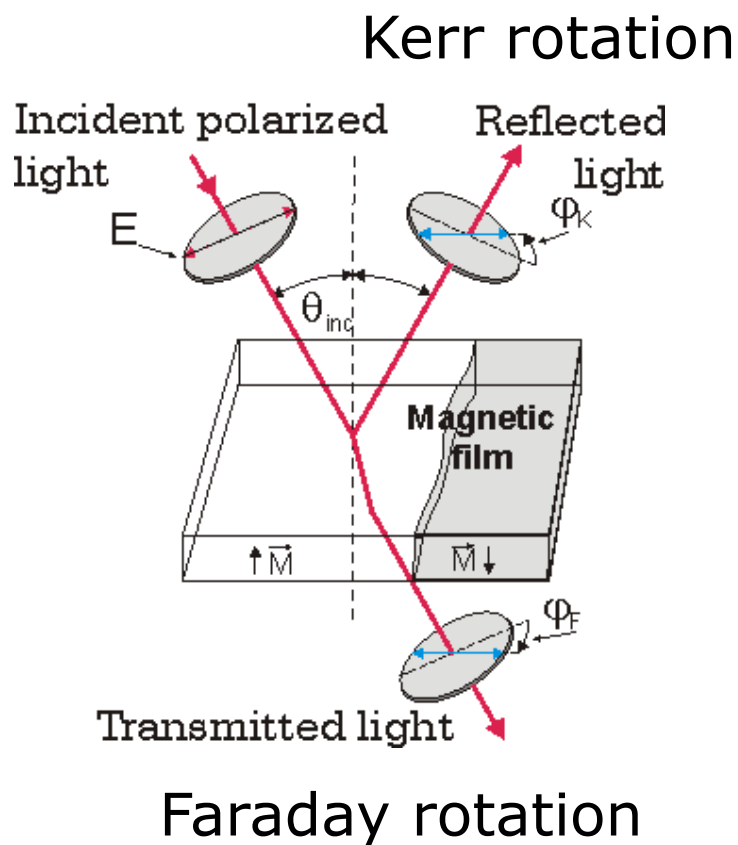
Boona, Myers, Heremans,
Energy and Env. Sci. (2014)



Motivation II: Can we make use of heat currents in information technology?

- Big picture problem: “How can we write magnetic information without resorting to magnetic fields, e.g., with spin currents?”
 - Rapid changes in magnetization and strong temperature gradients in magnetic materials should produce spin currents.
 - Magnitudes of the effects are only beginning to be understood.
- Why now and why pump-probe?
 - Leveraging rapidly advancing tools and growing knowledge base for heat transport in nanoscale metallic structures.
 - Create huge heat fluxes $100 \text{ GW m}^{-2} \text{ K}^{-1}$ and detect spin current in real time with 1 ps time resolution.

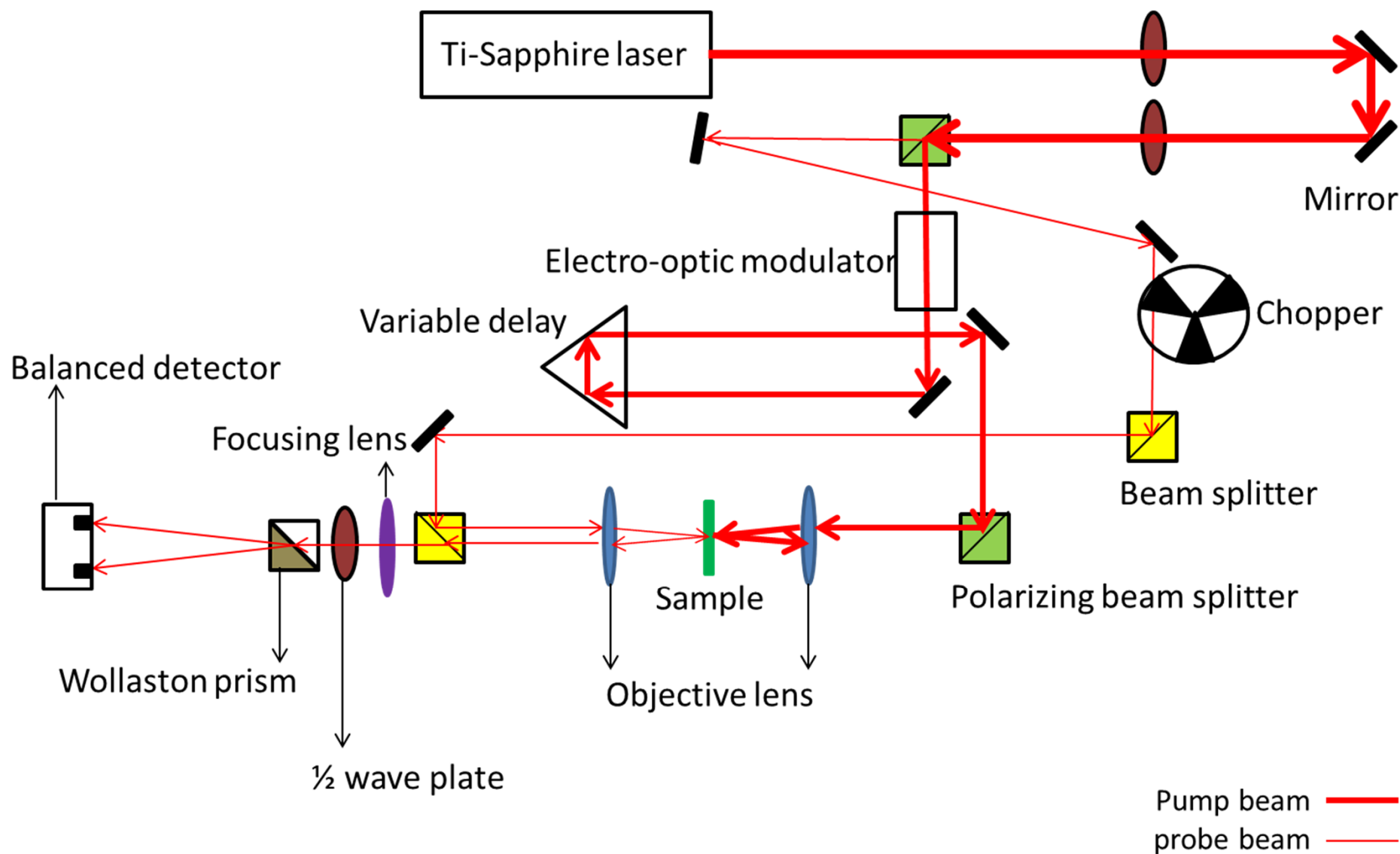
Time-resolved magneto-optic Kerr effect (TR-MOKE) to measure magnetization and spin accumulation



<http://labfiz.uwb.edu.pl>

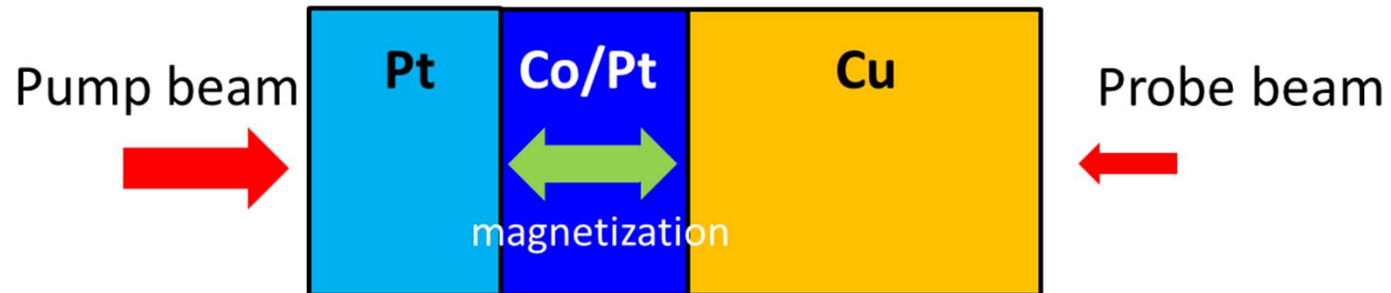
Körmann *et al.*, PRB (2011)

Time-resolved magneto-optic Kerr effect (TR-MOKE) to measure magnetization and spin accumulation

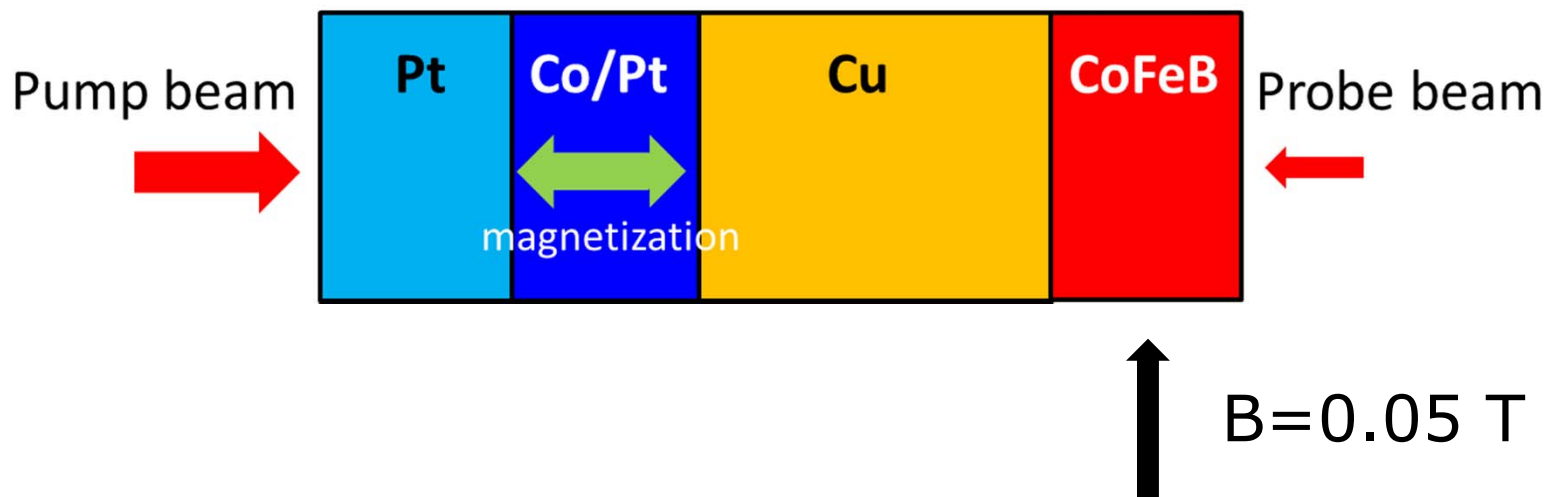


Two types of samples: i) for spin accumulation; and ii) for spin-transfer torque

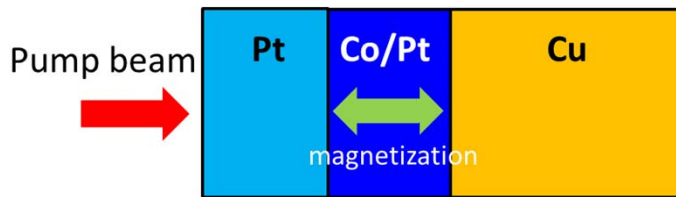
Sapphire/Pt(30)/[Co/Pt]_{xn}(6)/Cu(80)/MgO(10)/AlOx(5) (in nm)



Sapphire/Pt(30)/[Co/Pt]_{xn}(6)/Cu(10)/CoFeB(2)/MgO(10)/AlOx(5) (in nm)

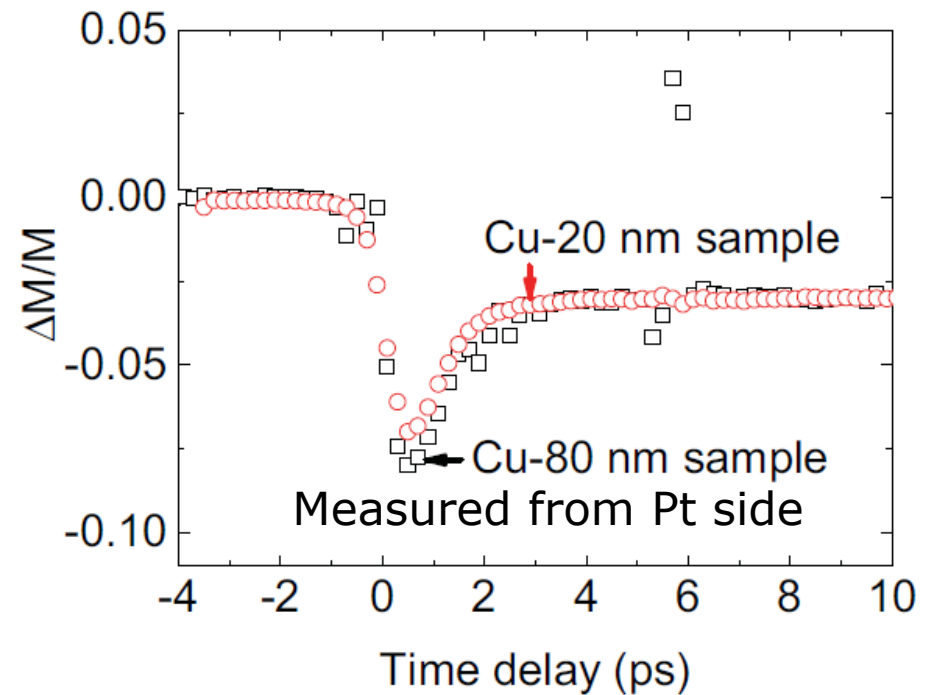
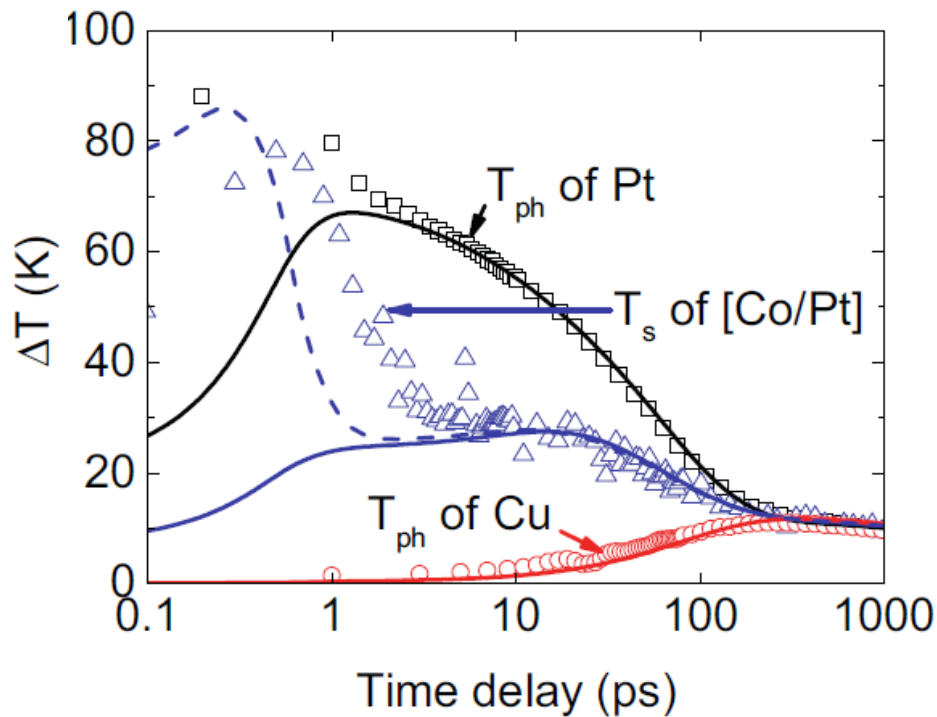


Pump Pt-side, probe either Pt-side or Cu side by either TDTR or TR-MOKE



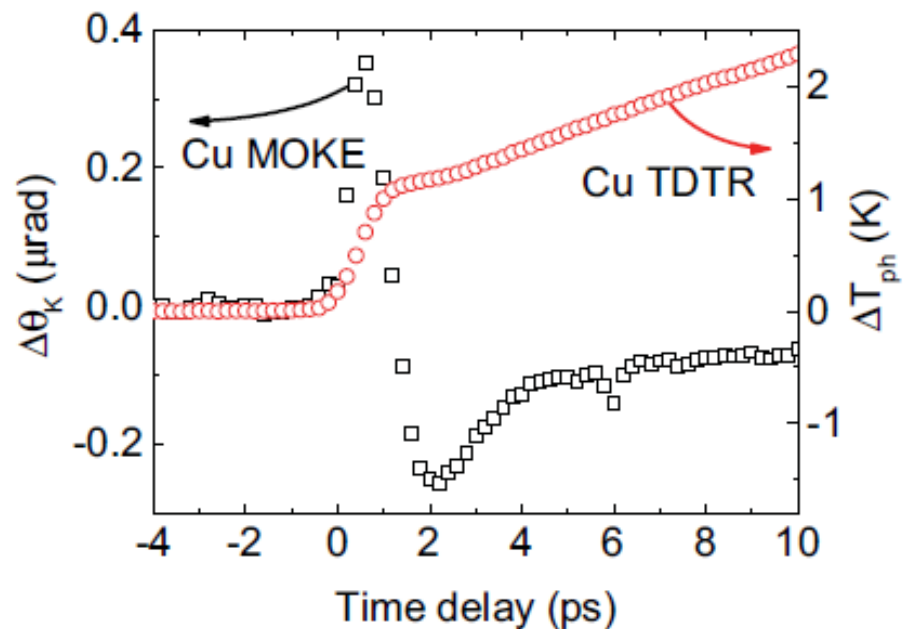
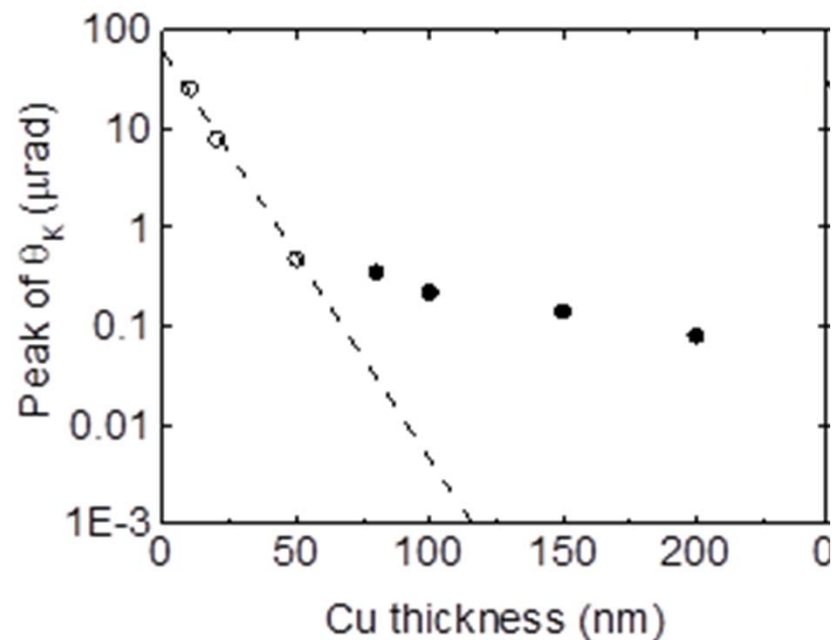
Normalized Kerr signal from Co/Pt is independent of Cu thickness

$$\left. \frac{\Delta M}{M} \right|_{\max} = -0.08 \pm 0.02 \text{ at } 0.5 \text{ ps}$$



Use thicker Cu layers to isolate contribution from spin-polarization in Cu

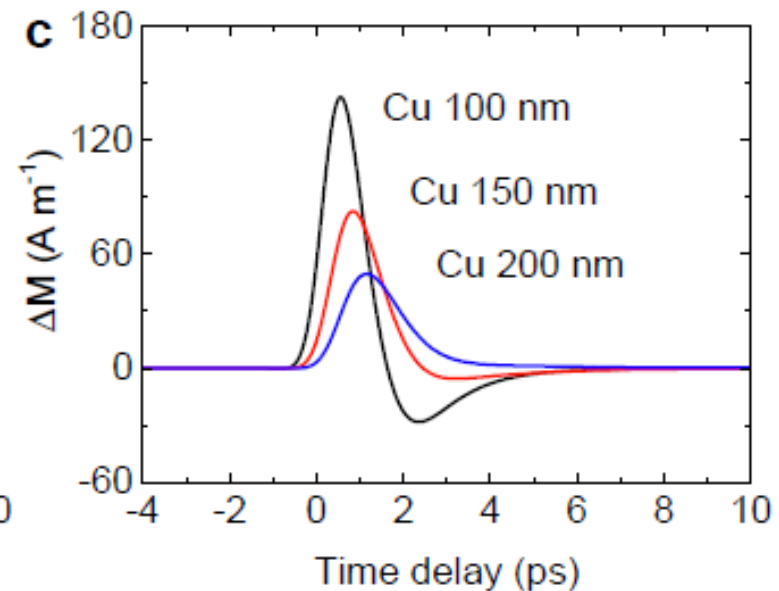
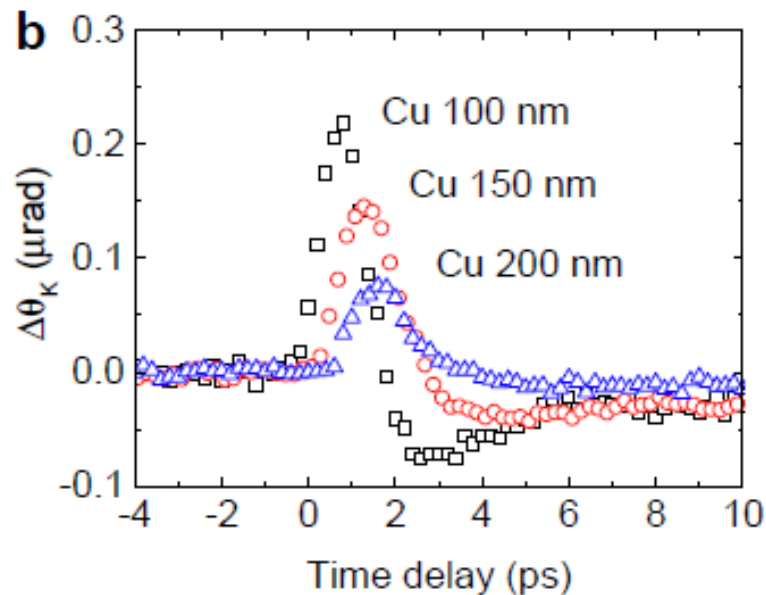
100 nm Cu



Comparison between experiment and spin diffusion model using spin generation = dM/dt

Measured Kerr signal on Cu side

Spin diffusion model



Spin diffusion model

$$\frac{\partial \mu_s}{\partial t} = D \frac{\partial^2 \mu_s}{\partial z^2} - \frac{\mu_s}{\tau_s}$$

spin generation rate per unit volume

$$G_s = -\frac{dM}{dt}$$

$\mu_s = \mu_{\uparrow} - \mu_{\downarrow}$ is the spin chemical potential

D is the spin diffusion constant

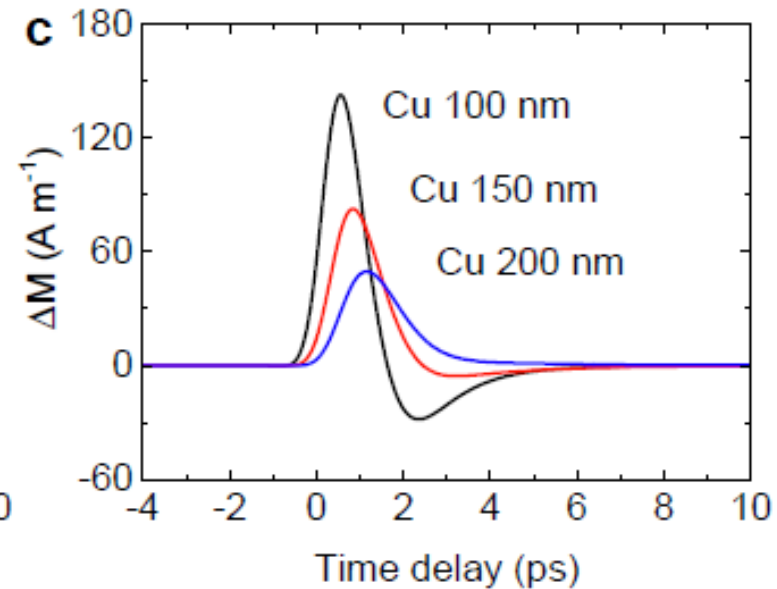
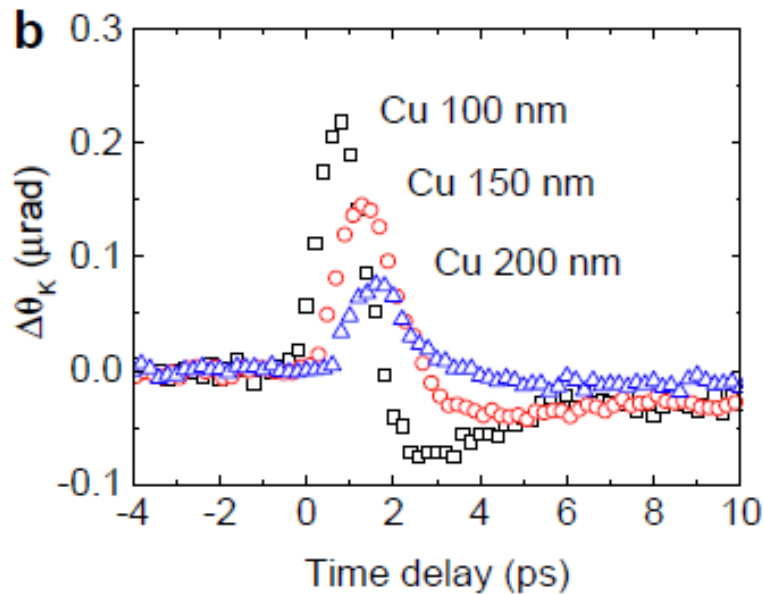
τ_s is the spin relaxation time.

	Pt	[Co/Pt]	Cu
D (nm ² /ps)	200	100	6500
τ_s (ps)	0.5	0.05	25
$(D\tau_s)^{1/2}$ (nm)	10	2.2	400

Comparison between experiment and spin diffusion model using spin generation = dM/dt

Measured Kerr signal on Cu side
 $E=36 \text{ J m}^{-2}$

Spin diffusion model
 $E=17 \text{ J m}^{-2}$



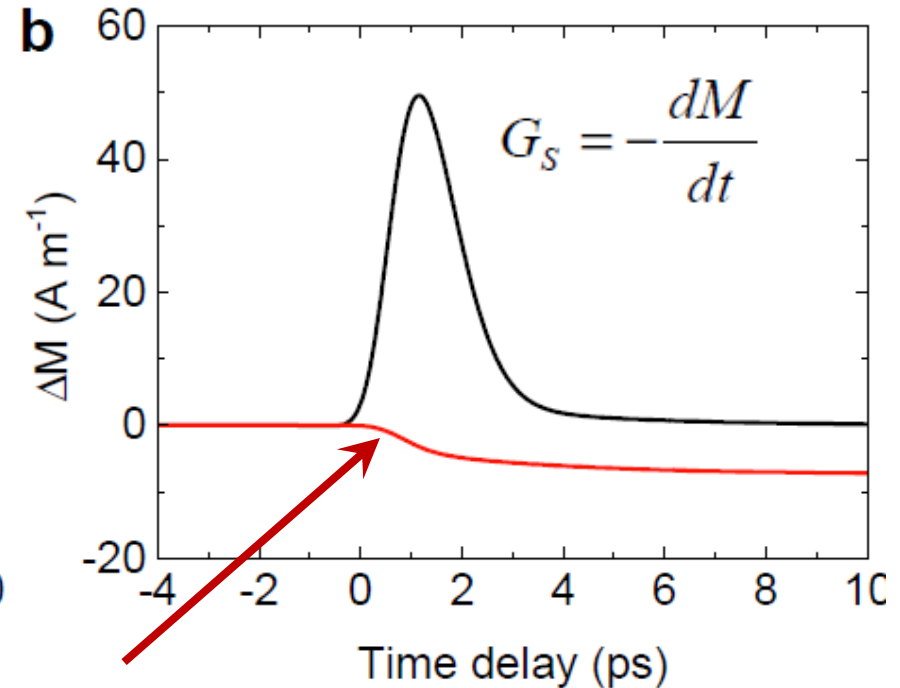
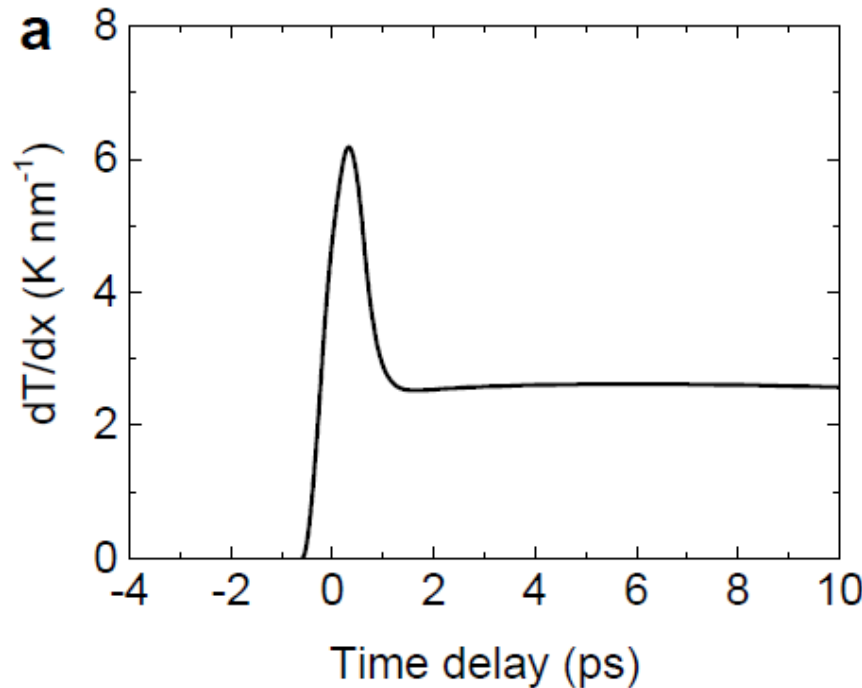
- No prior studies of how to convert Kerr rotation to spin accumulation.
- Working in progress to relate Kerr rotation quantitatively to spin accumulation in Cu and Au.

Temperature gradient also contributes to spin accumulation

$$J_s = -\frac{\mu_B}{e}(\sigma_{\uparrow}S_{\uparrow} - \sigma_{\downarrow}S_{\downarrow})\nabla T = -\frac{\mu_B}{e} \frac{\sigma_{\uparrow}S_{\uparrow} - \sigma_{\downarrow}S_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}} \sigma \nabla T$$

Temperature gradient in the Pt/Co layer from thermal modeling

Calculated spin accumulations



$$\frac{\sigma_{\uparrow}S_{\uparrow} - \sigma_{\downarrow}S_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}} \approx 5 \mu\text{V K}^{-1}$$

Spin diffusion modeling including spin-dependent Seebeck effect

- Spin chemical potential

$$\mu_s = \mu_{\uparrow} - \mu_{\downarrow} = A \exp\left(\frac{x}{l_s}\right) - A \exp\left(-\frac{x}{l_s}\right)$$

- Spin current for

$$|x| > l_s$$

$$\frac{d\mu_s}{dx} = \frac{A}{l_s} \exp\left(\frac{x}{l_s}\right) + \frac{A}{l_s} \exp\left(-\frac{x}{l_s}\right) \approx \frac{1}{l_s} \mu_s$$

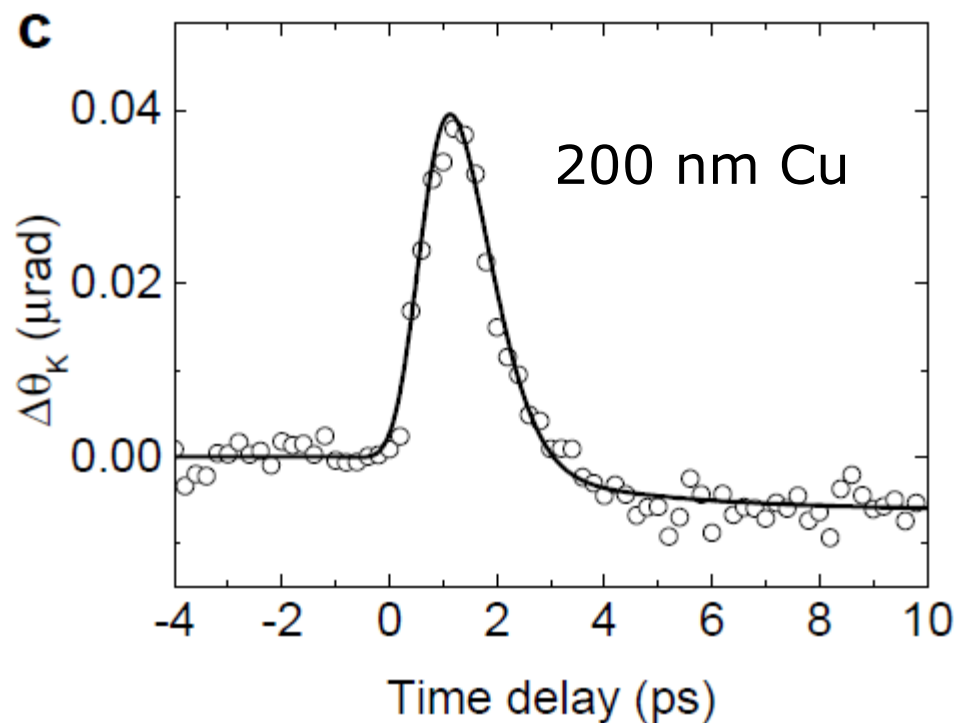
$$J_s = J_{\uparrow} - J_{\downarrow} = \Lambda_s \frac{d\mu_s}{dx} \quad \Lambda_s = \text{spin conductivity}$$

- Simple relationship between chemical potential and current

$$\mu_s = \frac{l_s}{\Lambda_s} J_s$$

Temperature gradient also contributes to spin accumulation

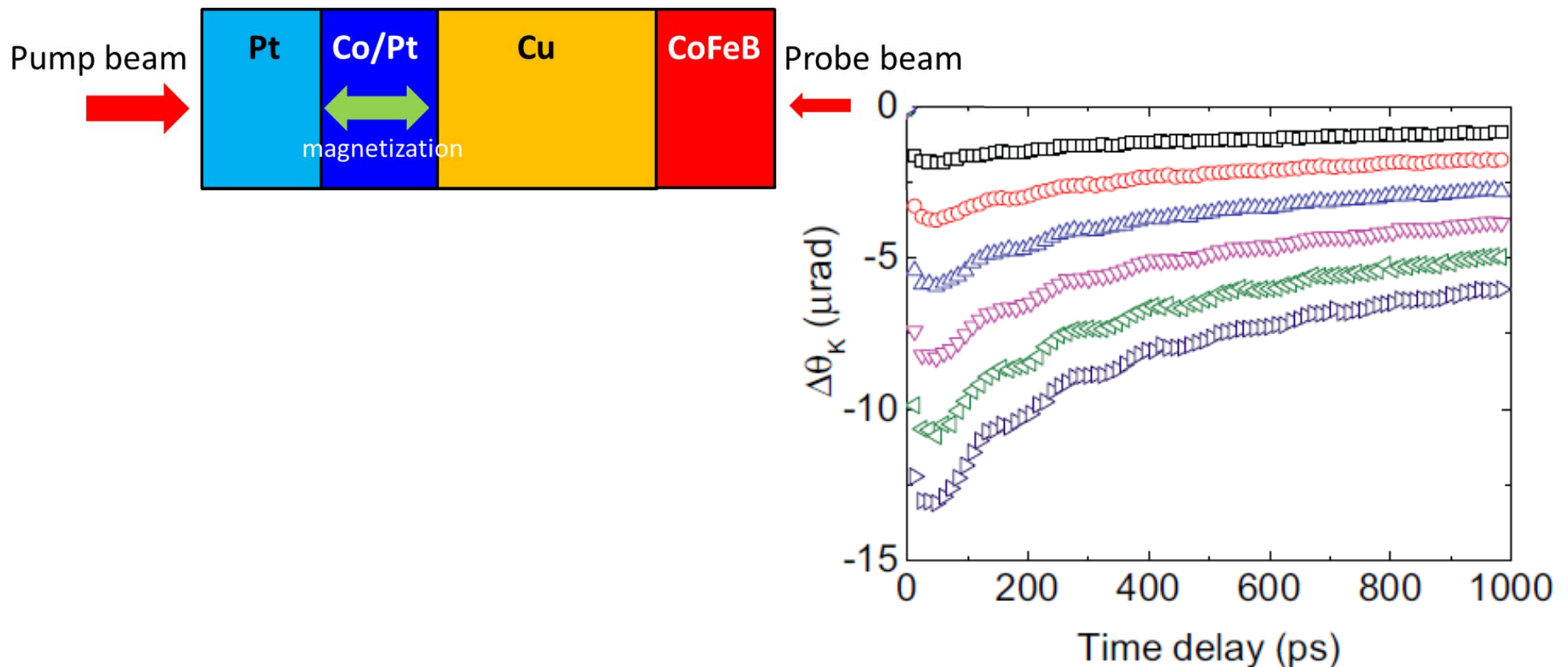
- More refined data with comparison to spin diffusion model including the spin-dependent Seebeck effect
- Comparison between model and data gives



$$\left. \frac{\Delta\theta_K}{\Delta M} \right|_{\text{Cu}} \approx 8.5 \times 10^{-10} \text{ rad m A}^{-1}$$

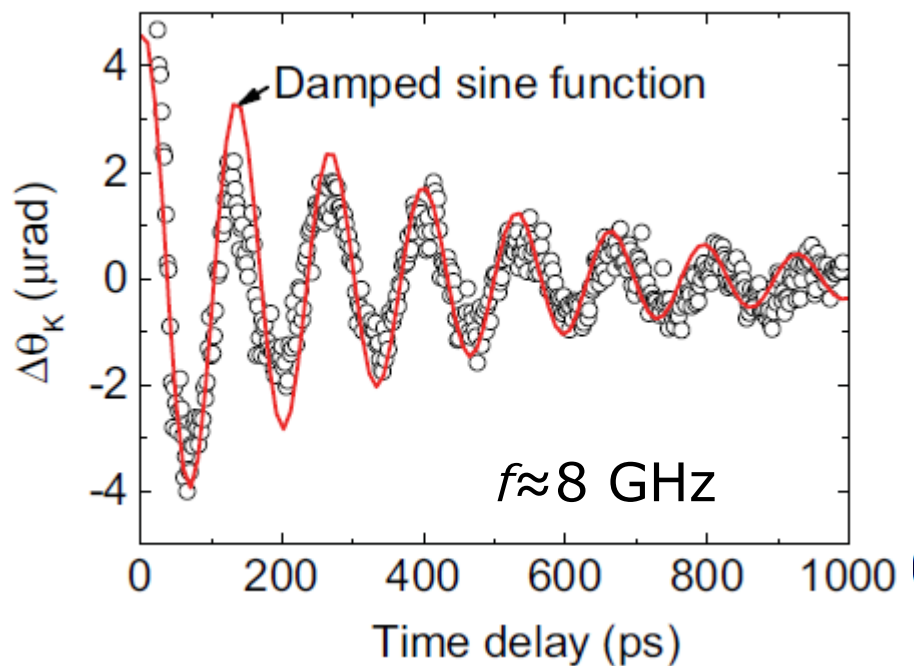
Use an in-plane magnetic layer of CoFeB as a “ballistic pendulum” for the spin current

- Spin current kicks magnetization of CoFeB out-of-plane (spin torque) and induces precession.
- Amplitude of the precession can be calibrated using Kerr rotation in a static field perpendicular field.

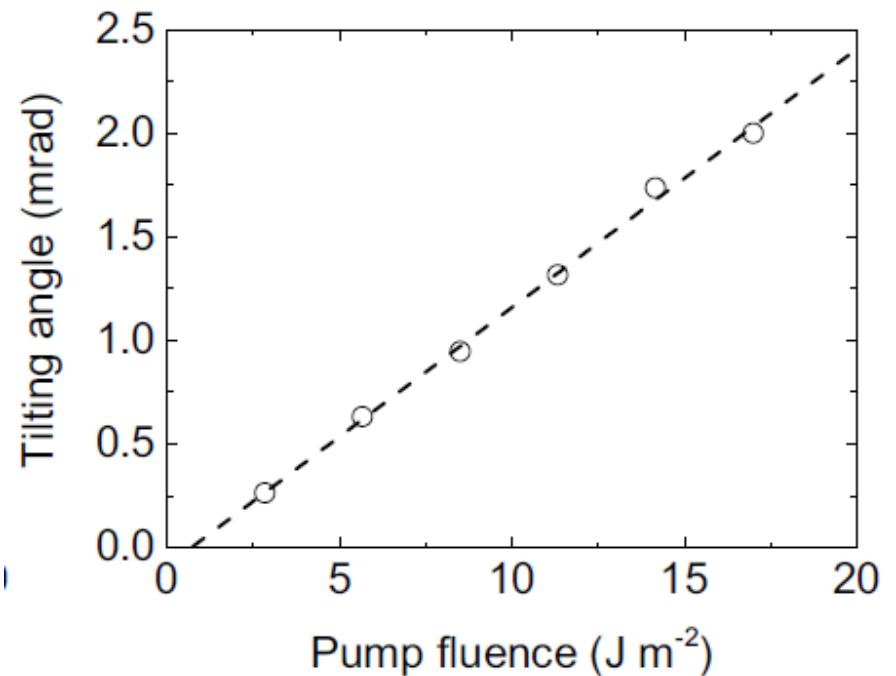


Precession frequency is well-described by Kittel equation

$$f = \frac{\gamma}{2\pi} \sqrt{\mu_0 H_x \left(\mu_0 H_x + \mu_0 M_S - \frac{2K_Z}{M_S} \right)},$$



Linear response



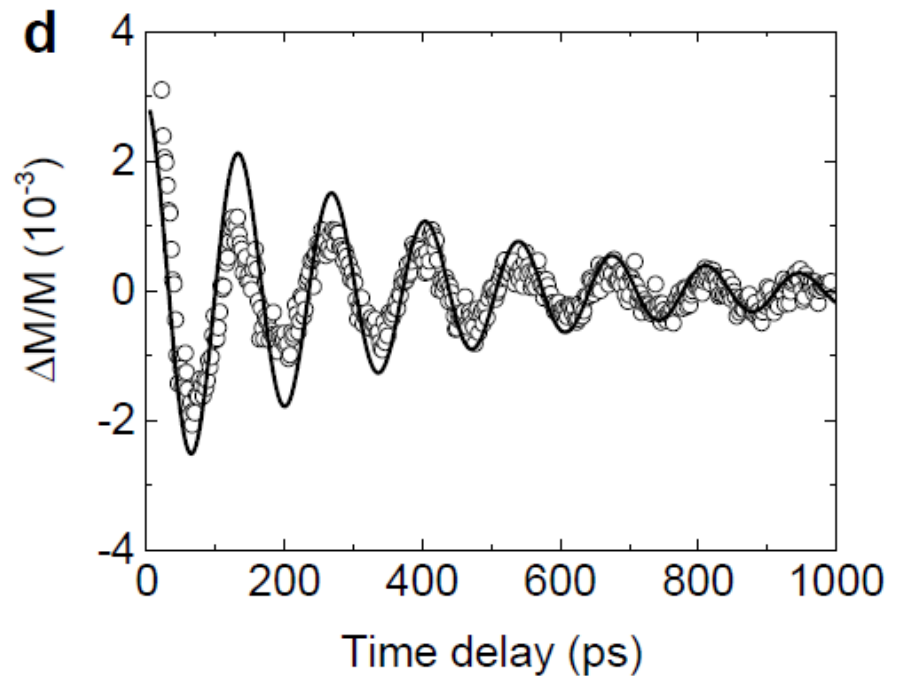
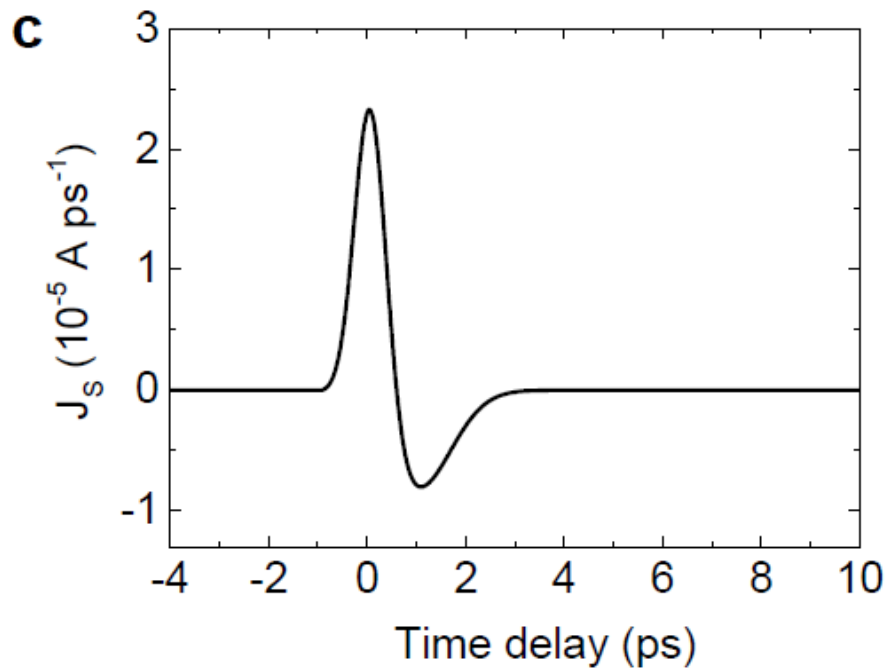
Combine spin diffusion model with magnetization dynamics

- Spin current has transverse polarization with respect to CoFeB magnetization, therefore, CoFeB is a perfect sink for spin (spin chemical potential is zero at Cu/CoFeB interface)
- Cu layer is thin, therefore, we need to include finite spin conductance at the [Co/Pt]/Cu and Cu/CoFeB interfaces
 - longitudinal spin conductance $\frac{G_{\uparrow} + G_{\downarrow}}{2e^2} \approx 0.4 \times 10^{15} \Omega^{-1} \text{ m}^{-2}$
 - transverse spin conductance $\frac{\text{Re}\{G_{\uparrow\downarrow}\}}{e^2} \approx 0.6 \times 10^{15} \Omega^{-1} \text{ m}^{-2}$

Good agreement between predicted and measured amplitude of spin precession

Landau-Lifshitz-Gilbert equation

$$\dot{m} = -\gamma m \times H_{\text{eff}} + \alpha m \times \dot{m} + \frac{J_s}{M_s h} m \times (m \times m_{\text{fixed}})$$



Summary

- Picosecond demagnetization of [Co/Pt] multilayer produces spin-currents that can exert a spin-transfer torque on an in-plane magnetic layer or produce spin accumulation in Cu
 - 6% of loss of demagnetization of [Co/Pt] magnetization is transferred to CoFeB layer
 - Increase efficiency with [Co/Pd] or [Co/Ni] with longer spin diffusion length?
- Coefficient for converting Kerr rotation to spin accumulation in Cu is 0.85 nm A^{-1}
 - Initial experiments on Au suggest that the detection sensitivity t is a factor of 5 larger than Cu
- Experiments and modeling give a spin-dependent Seebeck effect in [Co/Pt] of $\approx 5 \text{ } \mu\text{V K}^{-1}$
 - Will a tunnel barrier produce a larger effect?